## New techniques in dual-Doppler radar wind analysis

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## Goal of dual-Doppler wind analysis

Synthesize 3D fields of $u, v, w$ in small- and meso-scale phenomena from volume scans of radial wind data from two Doppler radars.


## Main ingredients of dual-Doppler wind analysis

Many different analysis frameworks (Cartesian vs co-plane, direct vs iterative, strong vs weak constraints), but most have same ingredients:

1. Radial wind observations $v_{\mathrm{r}}$ from two radars,
2. Smoothness constraint (explicit or pre/post processing filter),
3. Mass conservation (e.g., $\nabla \cdot[\bar{\rho}(z) \vec{u}]=0)$,
4. Impermeability condition ( $w=0$ at ground level and/or storm top).

## An exact theory for dual-Doppler wind analysis

Armijo (1969) derived the solution for a 3D velocity field $\vec{u}$ for which
(i) radial components of $\vec{u}$ agree with radial wind observations,

$$
\begin{align*}
& \vec{u} \cdot \hat{r}_{1}=v_{r 1},  \tag{4}\\
& \vec{u} \cdot \hat{r}_{2}=v_{r 2}, \tag{5}
\end{align*}
$$

(ii) anelastic mass conservation equation is satisfied,

$$
\begin{equation*}
\nabla \cdot[\bar{\rho}(z) \vec{u}]=0, \tag{6}
\end{equation*}
$$

(iii) impermeability condition is satisfied ( $w=0$ at ground level).

The $u, v, w$ fields satisfy (4)-(6). Eliminating $u$ and $v$ in favor of $w$ yields a 1st order partial differential equation for $w$. Get the exact analytical solution by integrating a forcing term along characteristics.

## Coplane coordinate system

The characteristics in the Armijo theory are circles in a cylindrical coordinate system whose central axis connects the radars (baseline). To get $w$ at any point (e.g., A), integrate the forcing term (data) along the circle passing through $\mathbf{A}$. Start at the ground (B) where $w=0$.

Well-posedness condition: a unique solution for $w$ exists at $\mathbf{A}$ if there are data at all points from $\mathbf{A}$ to $\mathbf{B}$. No solution exists if data are missing between $A$ and $B$.


## A Cartesian form of dual-Doppler wind analysis

Can bypass Armijo procedure, and solve (4)-(6) iteratively in Cartesian coordinates (e.g., Brandes 1977; Ray et al. 1980; Hildebrand \& Mueller 1985; Dowell \& Bluestein 1997).


However, the iterative procedure does not always converge. Dowell \& Shapiro (2003) derived a stability condition that showed that Armijo's "well-posedness" condition was relevant even in Cartesian coordinates.

## Ongoing challenges with dual-Doppler wind analysis: problems and some (partial) solutions

Even in cases where the analysis is well posed (either Armijo's Coplane analysis or the iterative Cartesian analysis), dual-Doppler analyses are still subject to a number of practical difficulties.

Problem 1: Biases in the divergence can quickly accumulate in the integration process and yield catastrophic errors in $w$.

Solution: Use radial wind data and mass conservation equation $\nabla \cdot\left[\rho_{0}(z) \vec{u}\right]=0$ as weak constraints (least squares error) in a variational procedure, e.g. 3DVAR or 4DVAR.

We will look at a 3DVAR procedure later.

Problem 2: Non-simultaneous data collection can result in phase (location) errors in key features such as gust fronts and vortices.

Solution: Use "advection correction." Invoke the frozen-turbulence hypothesis to shift data from both radars to a common analysis time.


## Frozen-turbulence hypothesis

Frozen-turbulence hypothesis: patterns translate (shift) without change in shape or intensity. In the case of the reflectivity field $Z$, this implies:

$$
\begin{equation*}
\frac{D Z}{D t}=0, \quad \text { or } \quad \frac{\partial Z}{\partial t}+U \frac{\partial Z}{\partial x}+V \frac{\partial Z}{\partial y}=0 \tag{7}
\end{equation*}
$$

where $U, V$ are pattern-translation components (not wind velocity components).

Many methods are available to estimate $U, V$ (e.g., Gal Chen 1982), however these generally treat $U$ and $V$ as constants over the whole grid.

We will consider a procedure to derive/use spatially variable $U, V$ fields in advection correction.

## Problem 3: Missing low-level data due to earth curvature, ground

 clutter, or non-zero elevation angle of lowest sweep.

Solution: Extrapolate data from the lowest sweep down to the ground.
Alternatively, use an additional constraint, e.g. a vorticity equation.
We will also look at this later.

## Spatially variable advection correction

We seek $U(x, y), V(x, y)$ and reflectivity $Z(x, y, t)$ fields on horizontal or constant elevation angle surfaces that minimize the cost function:

$$
\begin{equation*}
J \equiv \iiint\left[\alpha\left(\frac{\partial Z}{\partial t}+U \frac{\partial Z}{\partial x}+V \frac{\partial Z}{\partial y}\right)^{2}+\beta\left|\nabla_{h} U\right|^{2}+\beta\left|\nabla_{h} V\right|^{2}\right] d x d y d t \tag{8}
\end{equation*}
$$

with $Z$ imposed at two effective data times, $t=0$ and $t=T$.
$\beta$ is a smoothness coefficient; $\alpha$ is a data coverage function (= 0 or 1 ) that satisfies $\frac{\partial \alpha}{\partial t}+U \frac{\partial \alpha}{\partial x}+V \frac{\partial \alpha}{\partial y}=0$.

A similar $J$ underpins some single-Doppler velocity retrievals (Laroche \& Zawadzki 1995; Liou \& Luo 2001) and some precipitation nowcasting algorithms (Germann \& Zawadzki 2002).

## Euler-Lagrange equations

Two elliptic equations,

$$
\begin{align*}
& \beta T \frac{\partial^{2} U}{\partial x^{2}}+\beta T \frac{\partial^{2} U}{\partial y^{2}}=\int \alpha \frac{\partial Z}{\partial t} \frac{\partial Z}{\partial x} d t+U \int \alpha\left(\frac{\partial Z}{\partial x}\right)^{2} d t+V \int \alpha \frac{\partial Z}{\partial x} \frac{\partial Z}{\partial y} d t  \tag{9}\\
& \beta T \frac{\partial^{2} V}{\partial x^{2}}+\beta T \frac{\partial^{2} V}{\partial y^{2}}=\int \alpha \frac{\partial Z}{\partial t} \frac{\partial Z}{\partial y} d t+U \int \alpha \frac{\partial Z}{\partial x} \frac{\partial Z}{\partial y} d t+V \int \alpha\left(\frac{\partial Z}{\partial y}\right)^{2} d t \tag{10}
\end{align*}
$$

and one parabolic equation,

$$
\begin{array}{r}
\alpha\left(\frac{\partial}{\partial t}+U \frac{\partial}{\partial x}+V \frac{\partial}{\partial y}\right)^{2} Z+\alpha\left(\frac{\partial U}{\partial x}+\frac{\partial V}{\partial y}\right)\left(\frac{\partial Z}{\partial t}+U \frac{\partial Z}{\partial x}+V \frac{\partial Z}{\partial y}\right) \\
+\left(\frac{\partial \alpha}{\partial t}+U \frac{\partial \alpha}{\partial x}+V \frac{\partial \alpha}{\partial y}\right)\left(\frac{\partial Z}{\partial t}+U \frac{\partial Z}{\partial x}+V \frac{\partial Z}{\partial y}\right)=0 . \tag{11}
\end{array}
$$

## Analytical solution of the equation for $Z$

The characteristics of (11) are solutions of the trajectory equations: $\mathrm{D} x / \mathrm{D} t=U, \mathrm{D} y / \mathrm{D} t=V$. In characteristic coordinates, (11) becomes:

$$
\begin{equation*}
\left.\frac{D^{2} R}{D t^{2}}+\left(\frac{\partial U}{\partial x}+\frac{\partial V}{\partial y}\right) \right\rvert\, \frac{D R}{D t}=0 \tag{12}
\end{equation*}
$$

Integrate (13) twice with respect to time along trajectories. Evaluate constants of integration using data at the two input times. Solution is:

$$
\begin{equation*}
R(t)=R\left(t_{1}\right)+\left[R\left(t_{2}\right)-R\left(t_{1}\right)\right] \frac{I(t)}{I\left(t_{2}\right)} \tag{13}
\end{equation*}
$$

where

$$
\begin{equation*}
I(t) \equiv \int_{t_{1}}^{t} \exp \left[\left.-\int_{t_{1}}^{t^{\prime}}\left(\frac{\partial U}{\partial x}+\frac{\partial V}{\partial y}\right) \right\rvert\,\left(t^{\prime \prime}\right) d t^{\prime \prime}\right] d t^{\prime} \tag{14}
\end{equation*}
$$

## Analysis grid



## Combined analytical/numerical solution

Step 0: Construct CAPPIs at two data input times (2 volume scans)

Then, iterate between these steps:

Step 1: Solve the elliptic equations for $U$ and $V$ by relaxation.
Step 2: Calculate forward and backward trajectories running through all analysis points at all computational times.

Step 3: Interpolate $Z$ data to the end-points of each trajectory.
Step 4: Evaluate analytical solution for $Z$.

## Advection of reflectivity blobs in a solid body vortex



Advection-corrected $Z$ at middle time ( $t=\mathbf{3} \mathbf{~ m i n}$ )
Advection-corrected $Z$ is from a $\beta=100 \mathrm{dBZ}^{2}$ experiment.

## Retrieved $U$ in solid body vortex experiment



## Test case: Oklahoma supercell storm, 8 May 2003

Input data: Two scans of WSR-88D radar reflectivity (KTLX radar)



## Tests using 8 May 2003 TDWR data

Results with TDWR data were similar to those with WSR-88D data, but since TDWR data were available every $\sim 1 \mathrm{~min}$, could compare retrieved $Z$ with true $Z$. RMS error in $Z(\sim 4.5 \mathrm{dBZ})$ was less than RMS errors in $\mathbb{Z}$ obtained in any constant $U, V$ experiment:


## Use of the anelastic vertical vorticity equation in dual-Doppler wind analysis

Taking $\hat{k} \cdot(\nabla \times$ anelastic equations of motion) yields an equation for the evolution of the vertical vorticity $(\zeta=\partial v / \partial x-\partial u / \partial y)$ :

$$
\begin{equation*}
\left(\frac{\partial}{\partial t}+u \frac{\partial}{\partial x}+v \frac{\partial}{\partial y}+w \frac{\partial}{\partial z}\right) \zeta=\frac{\partial u \frac{\partial w}{\partial z}}{\partial y}-\frac{\partial v}{\partial z} \frac{\partial w}{\partial x}-\zeta\left(\frac{\partial u}{\partial x}+\frac{\partial v}{\partial y}\right) \tag{15}
\end{equation*}
$$

No baroclinic term in here (no $p$ or $\rho$ ). Baroclinicity is very important in convective storms, but the baroclinic vector is mostly horizontal.

Since (15) relates $w$ to $u$ and $v$, it can be used as a constraint in dual-Doppler wind analysis (Protat \& Zawadzki 2000; Protat et al. 2001; Mewes \& Shapiro 2002; Liu et al. 2005; Shapiro et al. 2009).

## Contending with unsteady term in vorticity equation

Method 0. Ignore the term

$$
\frac{\partial}{\partial t}\left(\frac{\partial v}{\partial x}-\frac{\partial u}{\partial y}\right)=0
$$

## Method 1. Impose frozen turbulence constraint

Impose frozen-turbulence constraint (say, with spatially variable $U, V$ ):

$$
\begin{equation*}
\frac{\partial}{\partial t}\left(\frac{\partial v}{\partial x}-\frac{\partial u}{\partial y}\right)=-U \frac{\partial}{\partial x}\left(\frac{\partial v}{\partial x}-\frac{\partial u}{\partial y}\right)-V \frac{\partial}{\partial y}\left(\frac{\partial v}{\partial x}-\frac{\partial u}{\partial y}\right) . \tag{16}
\end{equation*}
$$

Method 2. Impose frozen turbulence with intrinsic evolution
As air parcels translate, let their vorticity change linearly with time. Requires estimates of the vorticity fields at two time levels - get these fields from a simplified dual-Doppler analysis without $w$.

## 3DVAR analysis with vorticity equation constraint

Seek $u, v, w$ that minimize the sum of errors in the analysis constraints:

$$
\begin{align*}
J \equiv & \iiint \int\left(\alpha_{1} O_{1}^{2}+\alpha_{2} O_{2}^{2}\right) d r d \theta d \phi d t+  \tag{17}\\
& \iiint\left(\delta \varepsilon_{m}^{2}+\gamma \varepsilon_{v}^{2}+\beta_{1} S_{1}+\beta_{2} S_{2}+\beta_{3} S_{3}+\beta_{4} S_{4}\right) d x d y d z
\end{align*}
$$

$O_{1}, O_{2}$ : Differences between analyzed and observed $v_{\mathrm{r}}$ data.
$\varepsilon_{m}: \quad$ Residual in mass conservation equation.
$\varepsilon_{v}: \quad$ Residual in anelastic vertical vorticity equation.
$S_{1}-S_{4}: \quad$ Squared spatial derivatives of $u, v, w$ (smoothness terms).
$\beta_{1}-\beta_{4}$ : Smoothness weights.
$J$ is minimized with a conjugate-gradient algorithm.

## Test case: Oklahoma supercell storm, 8 May 2003



## Data denial experiments

## Control Run ("truth")

No vorticity equation constraint but all other constraints turned on. Radial wind data used throughout the analysis domain (as far down to the ground as possible; generally down to 100-200 m AGL)

## Data Denial Experiment 1: NOVORT

Radial winds thrown out for $z<1 \mathrm{~km}$. Otherwise, experiment is same as control run (no vorticity equation constraint).

## Data Denial Experiment 2: VORT

Radial winds thrown out for $z<1 \mathrm{~km}$. The vorticity equation constraint is turned on but with no provision for evolution.

## "True" $\boldsymbol{w}$ (m/s) at $z=1.75 \mathrm{~km}$ AGL



## Impact of vorticity constraint: <br> $w(\mathrm{~m} / \mathrm{s})$ at $z=1.75 \mathrm{~km}$ AGL



RMS error in $w$ versus height


## Tests using Advanced Regional Prediction System (ARPS) supercell storm data




## Data denial experiments

## Control Run ("truth")

w field output from ARPS run

## Data Denial Experiment 1: NOVORT

Radial winds thrown out for $z<1.5 \mathrm{~km}$. Data, mass conservation and smoothness constraints imposed. No vorticity equation constraint.

Data Denial Experiment 2: VORT
Radial winds thrown out for $z<1.5 \mathrm{~km}$. As in NOVORT but now the vorticity equation constraint is imposed. Spatially variable frozen turbulence applied in unsteady term but with no evolution.

## Data Denial Experiment 3: VORT+

Radial winds thrown out for $z<1.5 \mathrm{~km}$. As in VORT but now evolution is accounted for (crudely) in unsteady term.

$\boldsymbol{w}$ VORT ( 2 min scan time)

$w$ NOVORT

$w$ VORT+ ( 2 min scan time)



## 5 min volume scans



## 2 min volume scans



## 1 min volume scans



## 30 sec volume scans



## Focus on VORT+ results



## Future work

## 1. Advection Correction

Derive a spatially variable advection-correction procedure based on radial wind data, i.e., based on Euler-Lagrange equations arising from minimization of $D^{2}\left(r v_{r}\right) / D t^{2}=0$ subject to smoothness constraints.

## 2. Dual-Doppler wind analysis

Improve estimates of vorticity tendency by using improved spatially variable $U, V$ fields (see above).

Improve estimates of vorticity tendency by using rapid scan radar data (volume scans $\sim 1 \mathrm{~min}$ or less), e.g. from PAR radar, CASA radars, SMART-R radars, DOW radars.

